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CITATION:

Yamada, Osanobu. On the Spectrum of Dirac Operators with Potentials Diverging at Infinity (Spectral-scattering theory and related topics). 数理解析研究所講究録 1998, 1047: 126-133

ISSUE DATE:

1998-05

URL:

<http://hdl.handle.net/2433/62169>

RIGHT:

# On the Spectrum of Dirac Operators with Potentials Diverging at Infinity

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## 1 Results.

In this report we consider the spectrum of the Dirac operator

$$L = \sum_{j=1}^3 \alpha_j D_j + m(x) \beta + q(x) I_4 \quad \left( x \in \mathbf{R}^3, D_j = -i \frac{\partial}{\partial x_j} \right),$$

in the Hilbert space  $\mathcal{H} = [L^2(\mathbf{R}^3)]^4$ , where

$$\alpha_j = \begin{pmatrix} \mathbf{0} & \sigma_j \\ \sigma_j & \mathbf{0} \end{pmatrix} \quad (1 \leq j \leq 3), \quad \beta = \begin{pmatrix} I_2 & \mathbf{0} \\ \mathbf{0} & -I_2 \end{pmatrix}, \quad I_4 = \begin{pmatrix} I_2 & \mathbf{0} \\ \mathbf{0} & I_2 \end{pmatrix},$$

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad I_2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix},$$

and real valued functions  $m(x)$  and  $q(x)$  are assumed to be continuous in  $\mathbf{R}^3$  and satisfy

$$m(x) \longrightarrow +\infty \quad (\text{or} \quad -\infty) \quad \text{as} \quad r = |x| \longrightarrow \infty.$$

or

$$q(x) \longrightarrow +\infty \quad (\text{or} \quad -\infty) \quad \text{as} \quad r \longrightarrow \infty.$$

Let us denote the unique self-adjoint realization of  $L$  by  $H$ . We will show below the structure of the spectrum of the Dirac operator by dividing the problem into three cases ;

- (a)  $\limsup_{r \rightarrow \infty} \left| \frac{q(x)}{m(x)} \right| < 1,$
- (b)  $\limsup_{r \rightarrow \infty} \left| \frac{m(x)}{q(x)} \right| < 1,$
- (c)  $q(x) \equiv m(x) \quad (\text{or} \quad q(x) \equiv -m(x)).$

**Theorem A** (Yamada [9], Theorem 1). Assume (a). Let  $m, q \in C^1$  satisfy

$$(A.1) \quad m(x) \longrightarrow +\infty \text{ (or } -\infty) \text{ as } r \rightarrow \infty,$$

$$(A.2) \quad |\nabla m| = O(m(x)), \quad |\nabla q| = O(m(x)) \text{ as } r \rightarrow \infty.$$

Then, we have  $\sigma(H) = \sigma_d(H)$  (i.e., the set of discrete eigenvalues with finite multiplicity), which is unbounded at  $\pm\infty$  in  $\mathbf{R}$ .

**Theorem B** (Schmidt and Yamada [6], Theorem 1). Let  $m \in C^1$  and  $q \in C^0$  be spherically symmetric (i.e.,  $q = q(r)$ ,  $m = m(r)$ ), and satisfy (b) and

$$(B.1) \quad q(r) \longrightarrow +\infty \text{ (or } -\infty) \text{ as } r \rightarrow \infty,$$

$$(B.2) \quad \liminf_{r \rightarrow \infty} |m(r)| > 0,$$

$$(B.3) \quad \text{there exist a positive number } R_0 \text{ and two distinct real values } \lambda_1, \lambda_2 \text{ such that}$$

$$\frac{m}{q - \lambda_j} \in BV[R_0, \infty) \quad (j = 1, 2),$$

that is, they are of bounded variation in the interval  $[R_0, +\infty)$ .

$$(B.4) \quad \text{there exist a positive number } R_0 \text{ such that}$$

$$\frac{m'}{r m q} \in L^1(R_0, \infty).$$

Then, we have  $\sigma_{ac}(H) = \mathbf{R}$  and  $\sigma_s(H) = \emptyset$  where  $\sigma_{ac}(H)$  ( $\sigma_s(H)$ ) is the absolutely continuous (singular) spectrum of  $H$ .

**Remark 1.** In Theorem B, if  $m, q \in C^1$  satisfy (B.1), (B.2) and

$$\int_{R_0}^{\infty} \left( \left| \frac{m'}{q} \right| + \left| \frac{m q'}{q^2} \right| \right) dr < \infty$$

for some  $R_0 > 0$ , then (B.3) and (B.4) are satisfied.

**Theorem C** (Yamada [9], Theorem 2). Let  $m, q \in C^0$  satisfy

$$(C.1) \quad m(x) \equiv q(x) \longrightarrow +\infty \text{ as } r \rightarrow \infty,$$

Then, we have  $\sigma(H) \cap (0, +\infty) \subset \sigma_d(H)$ .

**Theorem C'** (Schmidt and Yamada [6], Theorem 2). Let  $q \in C^1$  be a spherically symmetric function. Assume (C.1) and

$$(C.2) \quad \text{there exists a positive number } R_0 \text{ such that}$$

$$\frac{q'}{q^{3/2}} \in BV[R_0, \infty) \cap L^2(R_0, \infty).$$

Then, we have

$$\sigma(H) \cap (-\infty, 0) \subset \sigma_{ac}(H), \text{ and } \sigma_s(H) \cap (-\infty, 0) = \emptyset.$$

**Remark 2.** In Theorem C', if  $q \in C^2$  satisfies

$$\int_{R_0}^{\infty} \left[ \frac{|q''|}{q^{3/2}} + \frac{(q')^2}{q^{5/2}} \right] dr < \infty,$$

then (C.2) is satisfied.

If  $m \equiv -q \rightarrow -\infty$ , then we have the similar result as in Theorem C'. On the other hand, if  $m \equiv q \rightarrow -\infty$ , or  $m \equiv -q \rightarrow +\infty$ , then see can see under the similar conditions that the negative spectrum is discrete, and the positive spectrum is absolutely continuous.

**Remark 3.** For the sake of simplicity we assumed the continuity of  $m(x)$  and  $q(x)$  in  $\mathbf{R}^3$ . It turns out that, if real valued functions  $m(x)$  and  $q(x)$  belong to  $L^2_{loc}(\mathbf{R}^3)$ , a symmetric operator  $L$  defined on  $C_0^\infty = [C_0^\infty(\mathbf{R}^3)]^4$  has at least one self-adjoint extension. For, the symmetric operator  $L$  is real with respect to a conjugation  $J$  such that

$$J u = \alpha_1 \alpha_3 \bar{u}.$$

## 2 Outline of the Proofs.

We sketch the proof of Theorem A, C, B, C', successively.

**The Proof of Theorem A.** It suffices to prove that  $(H - i)^{-1}$  is a compact operator on  $\mathcal{H}$ . Let  $\{f_n\}_{n=1,2,\dots}$  be a bounded sequence in  $\mathcal{H}$ , and set  $u_n = (H - i)^{-1} f_n$ . Then,  $u_n$  satisfies

$$(\alpha \cdot D) u_n + m(x) \beta u_n = [i - q(x)] u_n + f_n, \quad (1)$$

and, by operating  $(\alpha \cdot D)$  to (1),

$$\begin{aligned} (-\Delta + m^2 - q^2 + 1) u_n &= [\beta(\alpha \cdot D m) - 2i q - (\alpha \cdot D q)] u_n \\ &\quad + [(\alpha \cdot D) + (i - q + m \beta)] f_n \end{aligned}$$

in the distribution sense, where  $(\alpha \cdot D) = \sum_{j=1}^3 \alpha_j D_j$ . Using the assumptions we can find a positive constant  $C$  such that

$$\int_{\mathbf{R}^3} [|\nabla u_n|^2 + m^2(x) |u_n(x)|^2] dx \leq C \|f_n\|^2,$$

which and (A.1) imply the relative compactness of  $\{u_n\}$ .

**Remark 4.** We may adopt some local singularities of  $q(x)$  and  $m(x)$ . If we assume that there exist positive constants  $C$ ,  $R_0$  and  $\delta < 1$

$$\int_{|x| \leq R_0} |[m(x) \beta + q(x) I_4] u(x)|^2 dx \leq \delta \int_{|x| \leq R_0+1} |\nabla u|^2 dx + C \int_{|x| \leq R_0+1} |u(x)|^2 dx \quad (2)$$

for any  $u \in C_0^\infty$ , instead of the differentiability of  $m(x)$  and  $q(x)$  in  $|x| \leq R_0$ , then we obtain similarly from (1), (A.1) and (A.2) that

$$\int_{\mathbf{R}^3} |\nabla u_n|^2 dx + \int_{|x| \geq R_0} m^2(x) |u_n(x)|^2 dx \leq C \|f_n\|^2,$$

which implies also the relative compactness of  $\{u_n\}$ .

For example, if there exist positive constants  $R_0$  and  $\delta < 1$  such that

$$|m(x) \pm q(x)| \leq \frac{\delta}{2|x|} \quad (|x| \leq R_0),$$

then (2) holds in view of the well-known inequality

$$\int_{\mathbf{R}^3} \frac{|u(x)|^2}{|x|^2} dx \leq 4 \int_{\mathbf{R}^3} |\nabla u|^2 dx$$

for any  $u \in C_0^\infty$ .

**The Proof of Theorem C.** Let us show that any positive number  $\lambda$  does not belong to the essential spectrum  $\sigma_{ess}(H)$ . If otherwise, there exist a positive number  $\lambda$  and an orthonormal sequence  $\{u_n\}_{n=1,2,\dots}$  in  $\mathcal{H}$  such that

$$(H - \lambda) u_n \longrightarrow 0$$

in  $\mathcal{H}$ . Then, write

$$u_n = \begin{pmatrix} v_n \\ w_n \end{pmatrix}, \quad (H - \lambda) u_n = \begin{pmatrix} f_n \\ g_n \end{pmatrix} \longrightarrow \begin{pmatrix} 0 \\ 0 \end{pmatrix}.$$

Then, we have

$$\begin{cases} (\sigma \cdot D) w_n + 2q v_n - \lambda v_n = f_n \\ (\sigma \cdot D) v_n - \lambda w_n = g_n, \end{cases} \quad (3)$$

where  $(\sigma \cdot D) = \sum_{j=1}^3 \sigma_j D_j$ . Combining these identities we obtain

$$-\Delta v_n + 2\lambda q v_n = (\sigma \cdot D) g_n + \lambda f_n + \lambda^2 v_n. \quad (4)$$

which implies

$$\int_{\mathbf{R}^3} [|\nabla v_n|^2 + 2\lambda |q(x)| |v_n(x)|^2] dx \leq C (\|f_n\|^2 + \|g_n\|^2 + \|v_n\|^2). \quad (5)$$

where  $C$  is a positive constant independent of  $n = 1, 2, \dots$ . Thus, we can select a strongly convergent subsequence  $\{v_{n_j}\}_{j=1,2,\dots}$  in  $\mathbf{h} = [L^2(\mathbf{R}^3)]^2$ . Since  $u_n$  tends weakly to 0 in  $\mathcal{H}$ , we see

$$v_{n_j} \longrightarrow 0 \text{ in } \mathbf{h} \text{ as } j \rightarrow \infty,$$

which, (3) and (5) imply

$$w_{n_j} \longrightarrow 0 \text{ in } \mathbf{h} \text{ as } j \rightarrow \infty,$$

This fact contradicts to  $\|u_n\| = 1$  ( $n = 1, 2, \dots$ ).

**Remark 5.** In the proof of Theorem C, we may adopt a local singularity of  $q(x)$  in a ball, that is,

$$\int_{|x| \leq R_0} |q(x)|^2 dx < \infty,$$

instead of the continuity in the ball. For, if  $q(x)$  satisfies the above condition, then we have

$$\begin{aligned} \int_{|x| \leq R_0} |q(x)| |f(x)|^2 dx &\leq \delta \int_{|x| \leq R_0+1} |\nabla f|^2 dx + C_\delta \int_{|x| \leq R_0+1} |f(x)|^2 dx, \\ \int_{|x| \leq R_0} |q(x) f(x)|^2 dx &\leq \delta \int_{|x| \leq R_0+1} |\Delta f|^2 dx + C_\delta \int_{|x| \leq R_0+1} |f(x)|^2 dx \end{aligned} \quad (6)$$

for any  $f \in C_0^\infty$ , any positive number  $\delta < 1$ , and a positive constant  $C_\delta$ . Therefore, we can show (5) by means of (3) and (4). Thus, we can include Coulomb potentials in Theorem C without the restriction of the size of the constant.

**Remark 6.** If  $m(x) \equiv q(x)$  is an  $L_{loc}^2(\mathbf{R}^3)$  function, then the symmetric operator  $H_0$  with the domain  $D(H_0) = C_0^\infty$  such that  $H_0 u = L u$  ( $u \in D(H_0)$ ) is essentially self-adjoint. Indeed, we can see that the ranges of  $(H_0 \pm i)$  are dense in  $\mathcal{H}$ . If otherwise, we could take non-zero vectors  $v$  and  $w \in \mathbf{h}$  such that

$$\begin{cases} (\sigma \cdot D) w + 2q v = \eta v \\ (\sigma \cdot D) v = \eta w, \end{cases} \quad (7)$$

where  $\eta = i$  or  $-i$ , and

$$-\Delta v + 2\eta q v = -v \quad (8)$$

in the distribution sense. Then, we have  $|\nabla v| \in L^2(\mathbf{R}^3)$  by means of the latter of (7) and  $\Delta v \in \mathbf{h}$  in view of the assumption and (6), (8), which give

$$\int_{|x| \leq R} [|\nabla v|^2 + |v|^2] dx = \int_{|x|=R} \operatorname{Re} \left\langle \frac{\partial v}{\partial r}, \bar{v} \right\rangle dS. \quad (9)$$

Since the right hand side of (9) tends to 0 by a sequence  $\{R_n\}_{n=1,2,\dots}$  with  $R_n \rightarrow \infty$ , we have  $v = 0$ , which and the second identity of (7), gives  $w = 0$ . This is a contradiction.

**The Proof Theorem B.** If  $m = m(r)$  and  $q = q(r)$  are spherical symmetric functions, the spectral problem of  $H$  is reduces to the one of one-dimensional Dirac operators

$$l_k = -i \sigma_2 \frac{d}{dr} + m(r) \sigma_3 + q(r) I_2 + \frac{k}{r} \sigma_1 \quad (k = \pm 1, \pm 2, \dots)$$

in  $[L^2(0, \infty)]^2$  (see, e.g., Arai [1]).

The proof of the Theorem B is given on the line of the following Lemma 1 given by Behncke [3], Theorem 1.

**Lemma 1.** Let  $m, p$  and  $q$  be real valued functions and belong to  $L^1_{loc}(0, \infty)$ , and

$$l = -i\sigma_2 \frac{d}{dr} + m(r)\sigma_3 + q(r)I_2 + p(r)\sigma_1.$$

Let  $h_0$  be the minimal operator such that

$$h_0 u = l u, \quad u \in D(h_0) = \{u \in C_0^\infty(0, \infty) \mid l u \in \mathbf{h}\},$$

which is a densely defined symmetric operator in  $\mathbf{h}$  (see, e.g., Weidmann [8], Theorem 3.7). and let  $h$  be a self-adjoint extension of  $h_0$  (note that  $h_0$  is a real operator). Assume that  $I$  is an interval such that every solution of

$$l v = \lambda v \quad (\lambda \in I) \tag{10}$$

is bounded at infinity. Then, we have

$$I \subset \sigma_{ac}(h) \quad \text{and} \quad \sigma_s(h) \cap I = \emptyset.$$

The above lemma is closely related to Gilbert-Pearson [4] and Weidmann [7]. There is also a direct proof by Schmidt [5], Appendix.

In order to obtain the boundedness of  $v$  of (10) at infinity, we prepare the following Lemma 2.

**Lemma 2.** Let  $M, P$  and  $Q$  be real valued functions and belong to  $L^1_{loc}(0, \infty)$  such that

$$\lim_{r \rightarrow \infty} Q(r) = \infty \tag{11}$$

$$\limsup_{r \rightarrow \infty} \frac{\sqrt{M(r)^2 + P(r)^2}}{Q(r)} < 1, \tag{12}$$

and

$$\frac{\sqrt{M^2 + P^2}}{Q - \sqrt{M^2 + P^2}}, \quad \frac{M}{Q - \sqrt{M^2 + P^2}}, \quad \frac{P}{Q - \sqrt{M^2 + P^2}} \in BV[R_0, \infty) \tag{13}$$

for some  $R_0 > 0$ . Then, every solution of

$$-i\sigma_2 \frac{d}{dr} v + M\sigma_3 v + P\sigma_1 v + Qv = 0$$

is bounded at infinity. If  $P \equiv 0$ , the condition (13) may be read as ;

$$\frac{M}{Q - M} \in BV[R_0, \infty) .. \tag{14}$$

Theorem B is shown by seeing the boundedness of any solution  $v$  of

$$l_k v = -i \sigma_2 \frac{d}{dr} v + m(r) \sigma_3 + q(r) v + \frac{k}{r} \sigma_1 v = \lambda v \quad (\lambda < 0, \quad k = \pm 1, \pm 2, \dots)$$

at infinity. If we set

$$Q = q - \lambda, \quad M = m, \quad P = \frac{k}{r}$$

in Lemma 2, we can guarantee (11), (12) and (13) in Lemma 2 by means of the assumptions (B.1), (B.2), (B.3) and (B.4). Therefore, we get the boundedness of  $v$  at infinity, and the absolute continuity of the spectrum of  $H$  in view of Lemma 1.

**Remark 7.** If  $q(r)$  and  $m(r)$  are locally bounded in  $[0, \infty)$ , every  $l_k$  ( $k = \pm 1, \pm 2, \dots$ ) is of limit point type at 0 and  $\infty$ . If  $m(r)$  is locally bounded near 0 and  $q(r)$  satisfies

$$|q(r)| \leq \frac{\sqrt{3}}{2r}$$

near 0, then every  $l_k$  is of limit point type at 0. If  $m(r) = b/r$  ( $b$  is a real constant) and

$$|q(r)|^2 \leq \left[ \frac{3}{4} + b^2 \right] \frac{1}{r^2}$$

near 0, then every  $l_k$  is of limit point type at 0 (see, e.g., Arai [1] and Yamada [10], where are more general results).

**The Proof of Theorem C'.** We shall make use of the Gilbert–Pearson theory (Gilbert–pearson [4], Behncke [2]), showing that the differential equation

$$l_k v = \left( -i \sigma_2 \frac{d}{dr} + q(r) \sigma_2 + q(r) I_2 + \frac{k}{r} \sigma_1 \right) v = \lambda v \quad (\lambda < 0) \quad (15)$$

does not possess a subordinate solution at infinity, that is, any non-trivial solutions  $v$  and  $w$  of (15) for  $\lambda < 0$  satisfy

$$\liminf_{r \rightarrow \infty} \frac{\int_{R_0}^{\infty} |v(s)|^2 ds}{\int_{R_0}^{\infty} |w(s)|^2 ds} > 0 \quad (16)$$

for some  $R_0 > 0$ . To this end, for a solution  $v = {}^t(v_1, v_2)$  of (15), we set

$$\tilde{v} = \begin{pmatrix} \tilde{v}_1 \\ \tilde{v}_2 \end{pmatrix} = \begin{pmatrix} \sqrt[4]{(2q - \lambda)/(-\lambda)} v_1 \\ \sqrt[4]{(-\lambda)/(2q - \lambda)} v_2 \end{pmatrix}.$$

Then,  $\tilde{v}$  satisfies

$$\left( -i \sigma_2 \frac{d}{dr} + M \sigma_1 + Q I_2 \right) \tilde{v} = 0, \quad (17)$$

where

$$M(r) = \frac{-q'}{2(2q - \lambda)}, \quad Q(r) = \sqrt{(-\lambda)(2q - \lambda)}, \quad P \equiv 0.$$



Under the assumptions (C.1) and (C.2) we have that the above  $M$  and  $Q$  satisfy the conditions (11), (12) and (14) in Lemma 2 and, therefore, any solution  $\tilde{v}$  of (17) is bounded at infinity, which implies

$$C^{-1} \leq |\tilde{v}(r)|^2 \leq C \quad (r \geq R_0)$$

for some positive constants  $R_0$  and  $C > 1$ . Thus, we have

$$0 < \liminf_{r \rightarrow \infty} \frac{|v(r)|^2}{\sqrt{2q(r) - \lambda}} \leq \limsup_{r \rightarrow \infty} \frac{|v(r)|^2}{\sqrt{2q(r) - \lambda}} < \infty.$$

The same estimate holds for  $w$ , which yields (16).

### References.

- [1] Arai, M., On essential selfadjointness, distinguished selfadjoint extension and essential spectrum of Dirac operators with matrix valued potentials, Publ. RIMS, Kyoto Univ., **19** (1983), 33–57.
- [2] Behncke, H., Absolute continuity of Hamiltonians with von Neumann Wigner potentials, Proc. Amer. Math. Soc., **111** (1991), 373–384.
- [3] Behncke, H., Absolute continuity of Hamiltonians with von Neumann Wigner potentials II, Manuscripta Mth., **71** (1991), 163–181.
- [4] Gilbert, D.J. and Pearson D.B., On subordinacy and analysis of the spectrum of one-dimensional Schrödinger operators, J. Math. Anal. Appl., **128** (1987), 30–56.
- [5] Schmidt, K.M., Absolutely continuous spectrum of Dirac systems with potentials infinite at infinity, Math. Proc. Camb. Phil. Soc. (to appear).
- [6] Schmidt, K.M. and Yamada, O., Spherically symmetric Dirac operators with variable mass and potentials infinite at infinity (preprint).
- [7] Weidmann, J., Oszillationsmethoden für Systems gewöhnlicher Differentialgleichungen, Math. Z., **119** (1971), 349–373.
- [8] Weidmann, J., *Spectral Theory of Ordinary Differential operators*, Lecture Notes in Mathematics **1258** (1987), Springer-Verlag.
- [9] Yamada, O., On the spectrum of Dirac operators with the unbounded potential at infinity, Hokkaido Math. J., **26** (1997), 439–449.
- [10] Yamada, O., A remark on the essential self-adjointness of Dirac operators, Proc. Japan Acad., Ser. A., **62** (1986), 327–330.